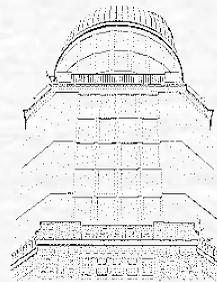


On the Relativistic Theory of Earth Rotation

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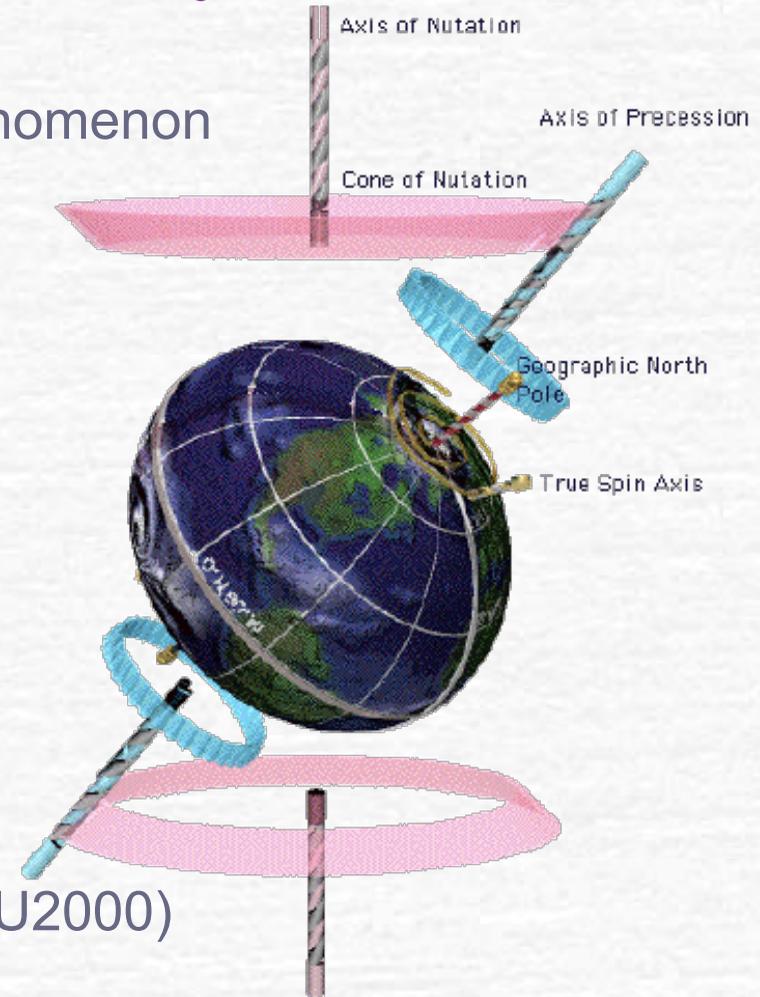
Relativity and Earth rotation: why to bother?

- Earth rotation is the only astronomical phenomenon
 - which is **observed with a high accuracy**

and

- which has **no widely-used consistent relativistic model**

- Modern theories of precession/nutation (IAU2000) are based on purely Newtonian theories with geodetic precession and nutation added in an inconsistent way
- Modern theories of rigid Earth nutation are intended to attain formal accuracy of $1 \mu\text{as}$ (expected relativistic effects are much larger)



How to model?

- Early attempts (- 1986)

- One single reference system **BCRS** for both translational motion of solar system and for rotational motion of all the bodies...
- The results were clearly physically inadequate coming from bad choice of coordinates:

E.g. spurious annual variations in LOD with an amplitude of **75 μ s...**

Reason: “**bad**” coordinates that provide
no analogy of Newtonian tidal forces at the post-Newtonian level

“Better” coordinates are clearly needed: **GCRS**

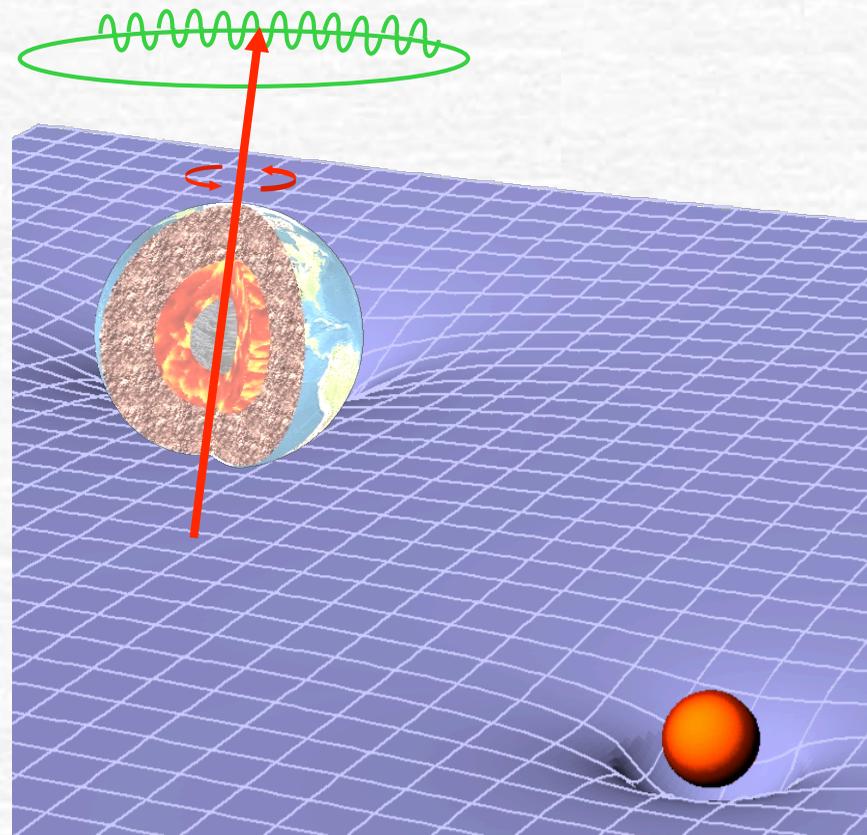
How to model?

- More sophisticated way (1986 -)
 - A physically adequate local **GCRS**
 - Still some coordinates, but chosen in such a way that the influence of external gravitational fields is as small as possible:

full analogy of Newtonian tidal forces at the post-Newtonian level

Main goal of the project

- Derivation of a new consistent and improved precession/nutation series for a **rigidly rotating** multipole model of the Earth in the post-Newtonian approximation of general relativity
- using post-Newtonian definitions of :
 - **potential coefficients**
 - **moment of inertia tensor**
- dynamical equations in the GCRS
- correct relativistic time scales
- rigorous treatment of the geodetic precession and nutation



Equations of rotational motion in the GCRS

- Post-Newtonian equations of rotational motion in the GCRS (Damour, Soffel, Xu, 1993, Klioner, Soffel et al 1996-)

$$\frac{d}{dT_{CG}} \left(C^{ab} \boldsymbol{\omega}^b \right) = \sum_{l=1}^{\infty} \frac{1}{l!} \boldsymbol{\varepsilon}_{abc} M_{bL} G_{cL} + L^a(\mathbf{C}, \boldsymbol{\omega}, \boldsymbol{\Omega}_{\text{iner}}) + \dots$$

The last terms is the Coriolis torque from the relativistic precessions:

$$\boldsymbol{\Omega}_{\text{iner}}^a = -\frac{3}{2c^2} \boldsymbol{\varepsilon}_{aij} v_E^i \frac{\partial}{\partial x^j} w_{\text{ext}}(\mathbf{x}_E) + \frac{2}{c^2} \boldsymbol{\varepsilon}_{aij} \frac{\partial}{\partial x^j} w_{\text{ext}}^i(\mathbf{x}_E) - \frac{1}{2c^2} \boldsymbol{\varepsilon}_{aib} v_E^i G_b$$

geodetic precession

Lense-Thirring precession

Thomas precession (negligible)

$$w_{\text{ext}}(\mathbf{x}) = \sum_{A \neq E} \frac{GM_A}{|\mathbf{x} - \mathbf{x}_A|}, \quad w_{\text{ext}}^i(\mathbf{x}) = \sum_{A \neq E} \frac{GM_A}{|\mathbf{x} - \mathbf{x}_A|} v_A^i$$

Rigidly rotating multipoles in the GCRS

- Klioner, Soffel, Xu, Wu, 2001 (based on many previous results):

- Post-Newtonian equations of rotational motion in the GCRS

$$\frac{d}{dTCG} \left(C^{ab} \omega^b \right) = \sum_{l=1}^{\infty} \frac{1}{l!} \varepsilon_{abc} M_{bL} G_{cL} + L^a (\mathbf{C}, \omega, \mathbf{\Omega}_{\text{iner}}) + \dots$$

- Rigidly rotating multipoles:

several assumptions on the multipole moments and
the tensor of inertia

$$C^{ab} = P^{ac} P^{bd} \bar{C}^{ab}, \quad \bar{C}^{ab} = \text{const},$$

$$M_{a_1 a_2 \dots a_l} = P^{a_1 b_1} P^{a_2 b_2} \dots P^{a_l b_l} \bar{M}_{b_1 b_2 \dots b_l}, \quad \bar{M}_{b_1 b_2 \dots b_l} = \text{const}, \quad l \geq 2,$$

$$\omega^a = \frac{1}{2} \varepsilon_{abc} P^{db} \frac{d}{dTCG} P^{dc}$$

$P^{ab}(TCG)$ is an orthogonal matrix defining the orientation of the ITRS in GCRS

Numerical code: an overview

- Fortran 95, about 20000 lines
- careful coding to avoid excessive numerical errors
- two numerical integrators: ODEX and ABM with dense output
- automatic accuracy check: forth and back integrations
- any available arithmetic: 64 bit, 80 bit, 128 bit
- extended-precision arithmetic for precision-critical operations (switchable)
- the STF code has been automatically generated by *Mathematica*
- baseline:

ODEX with 80 bits on Intel architecture gives errors $<0.001 \mu\text{s}$ for 150 years

- in the Newtonian limit reproduces SMART within the errors of the latter
- performance:

Newtonian case:	2.2 sec per yr
all the relativity on:	8.8 sec per yr

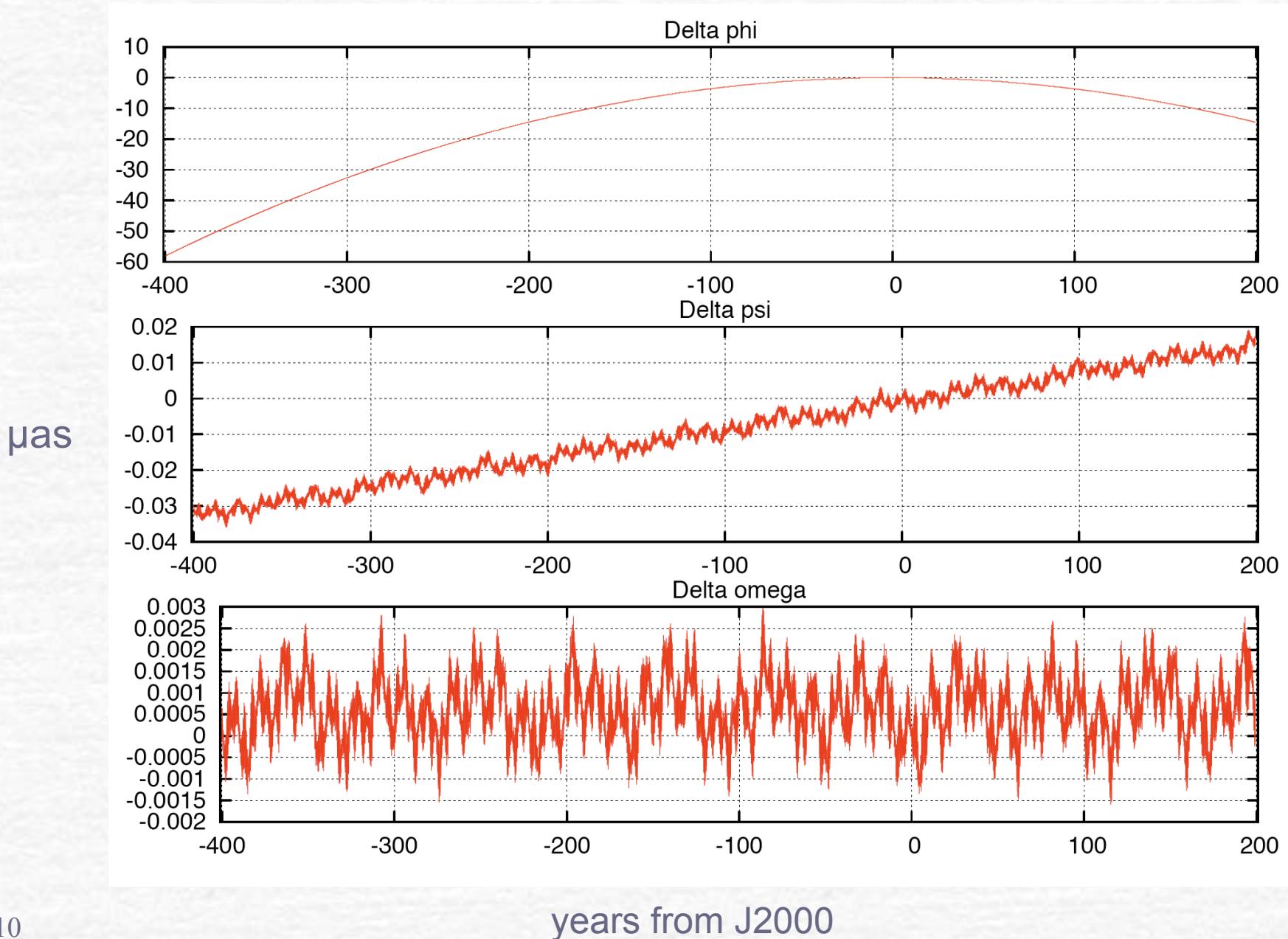
Long-term numerical integrations

A first step to a relativistic theory of precession ...

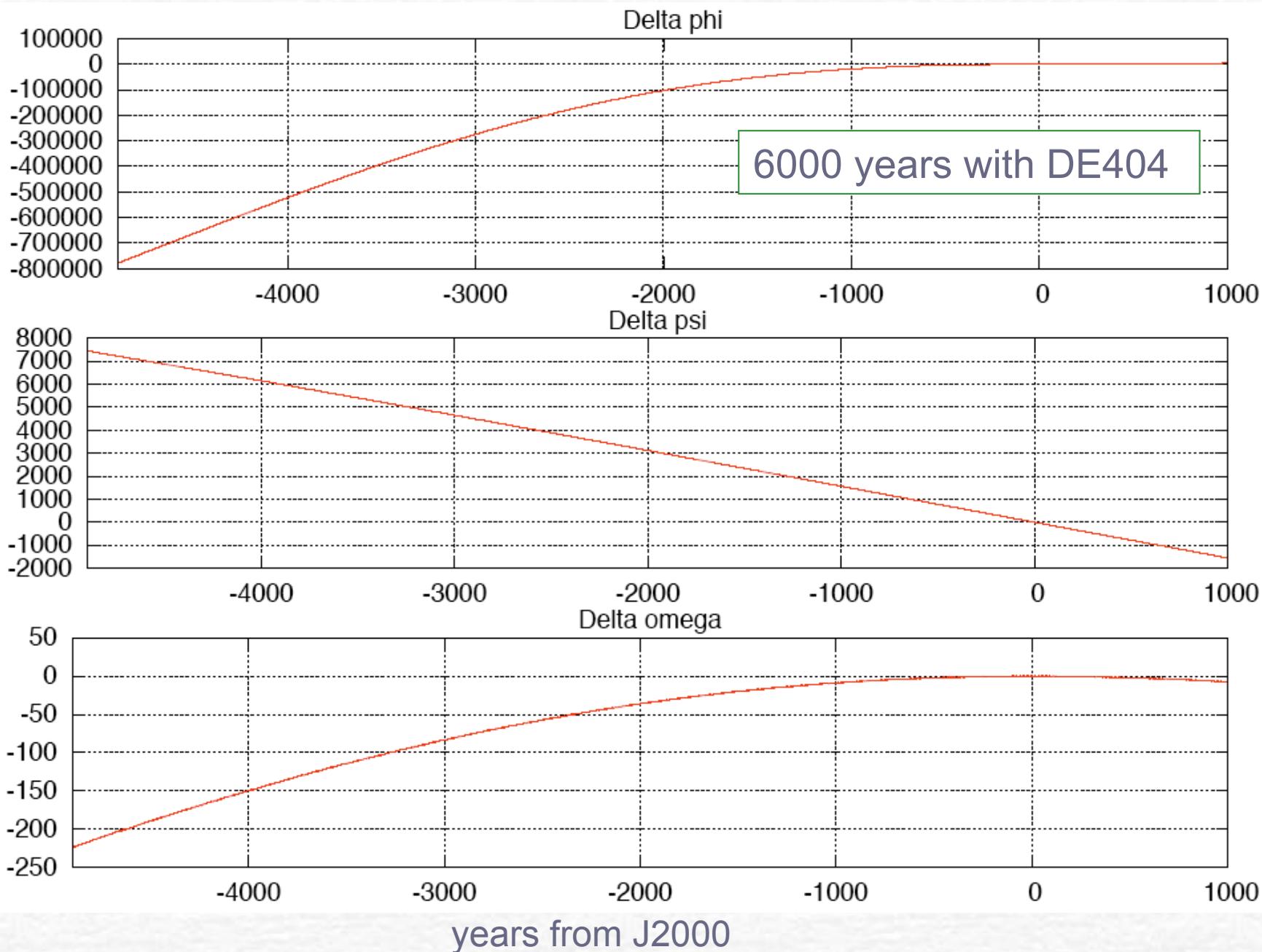
A long solar system ephemeris is needed:

DE404 is used to check the situation: 6000 yr

Noise from the downgrade: DE404 vs DE403



Effects of the post-Newtonian torque



Effects of the post-Newtonian torque

6000 years with DE404:

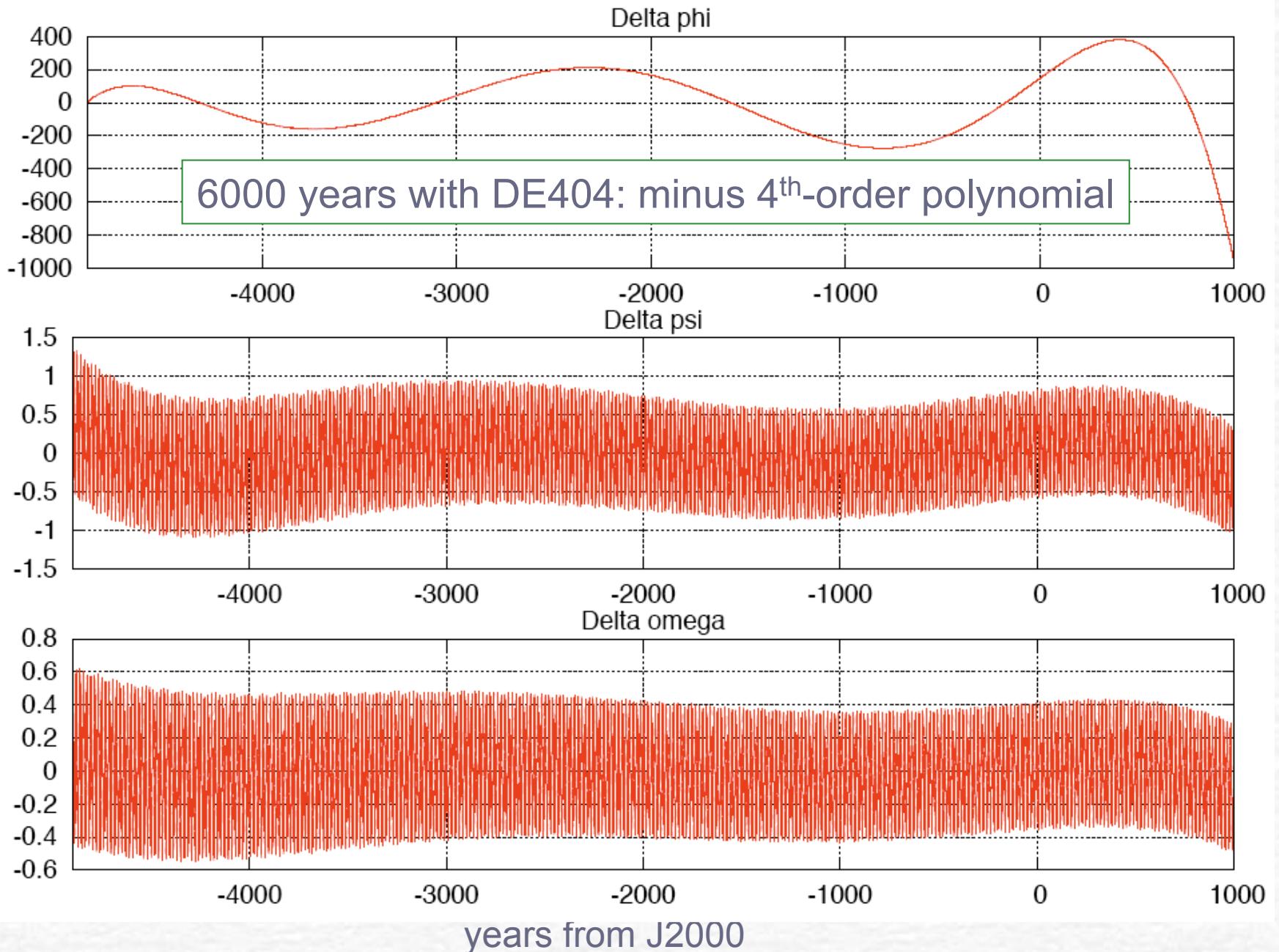
$$\Delta\varphi = -146.67 + 640.60t - 7921.70t^2 \\ + 11375.50t^3 + 1308.23t^4$$

$$\Delta\psi = -0.61 - 1560.31t + 3.22t^2 + 2.13t^3 - 0.06t^4$$

$$\Delta\omega = -0.23 + 0.014t - 7.99t^2 + 0.64t^3 + 0.08t^4$$

in μas , t is in thousand years

Effects of the post-Newtonian torque

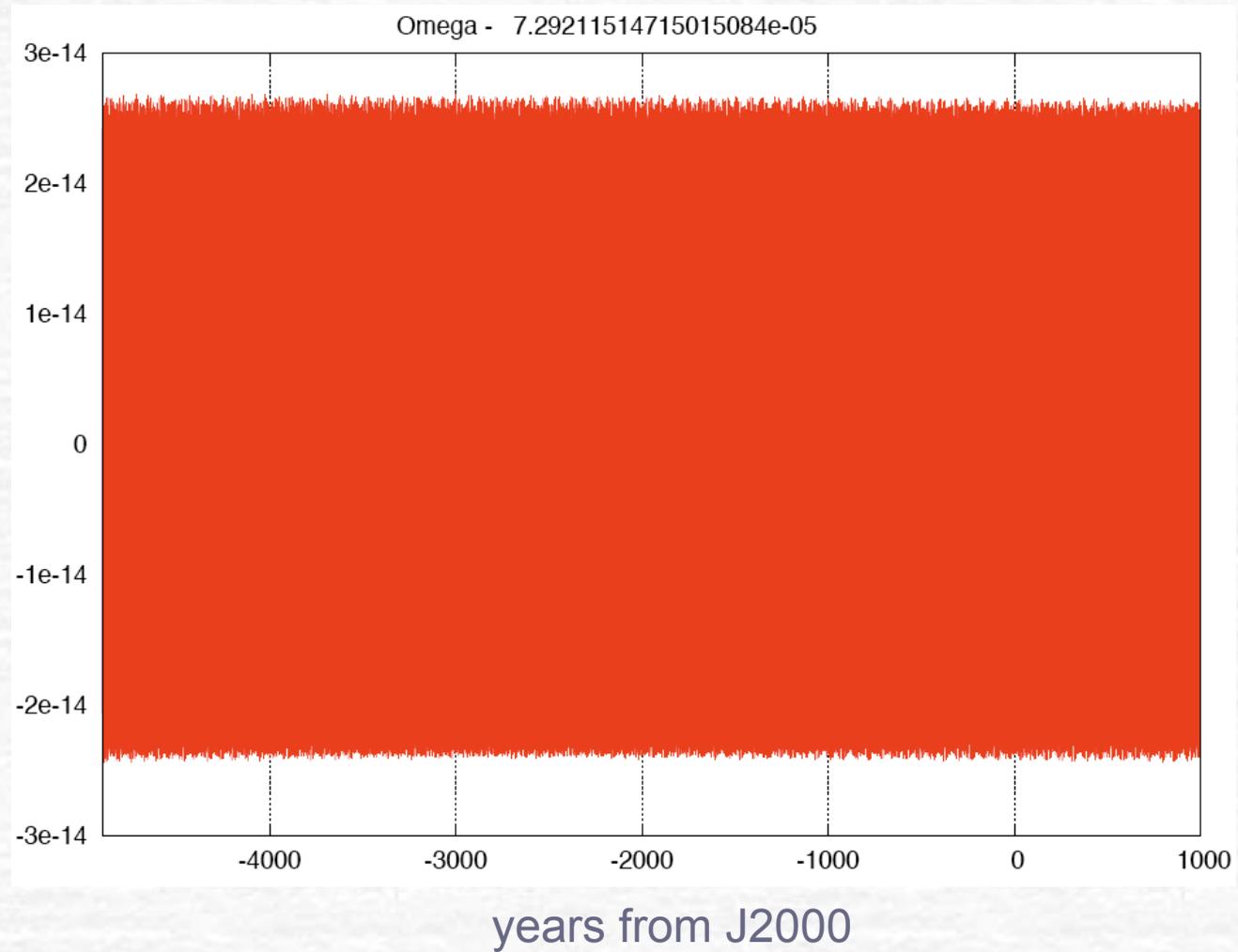


Effects on the LOD

6000 years with DE404:

$$\dot{\varphi} + \dot{\psi} \cos \omega = \Omega$$

rad/s

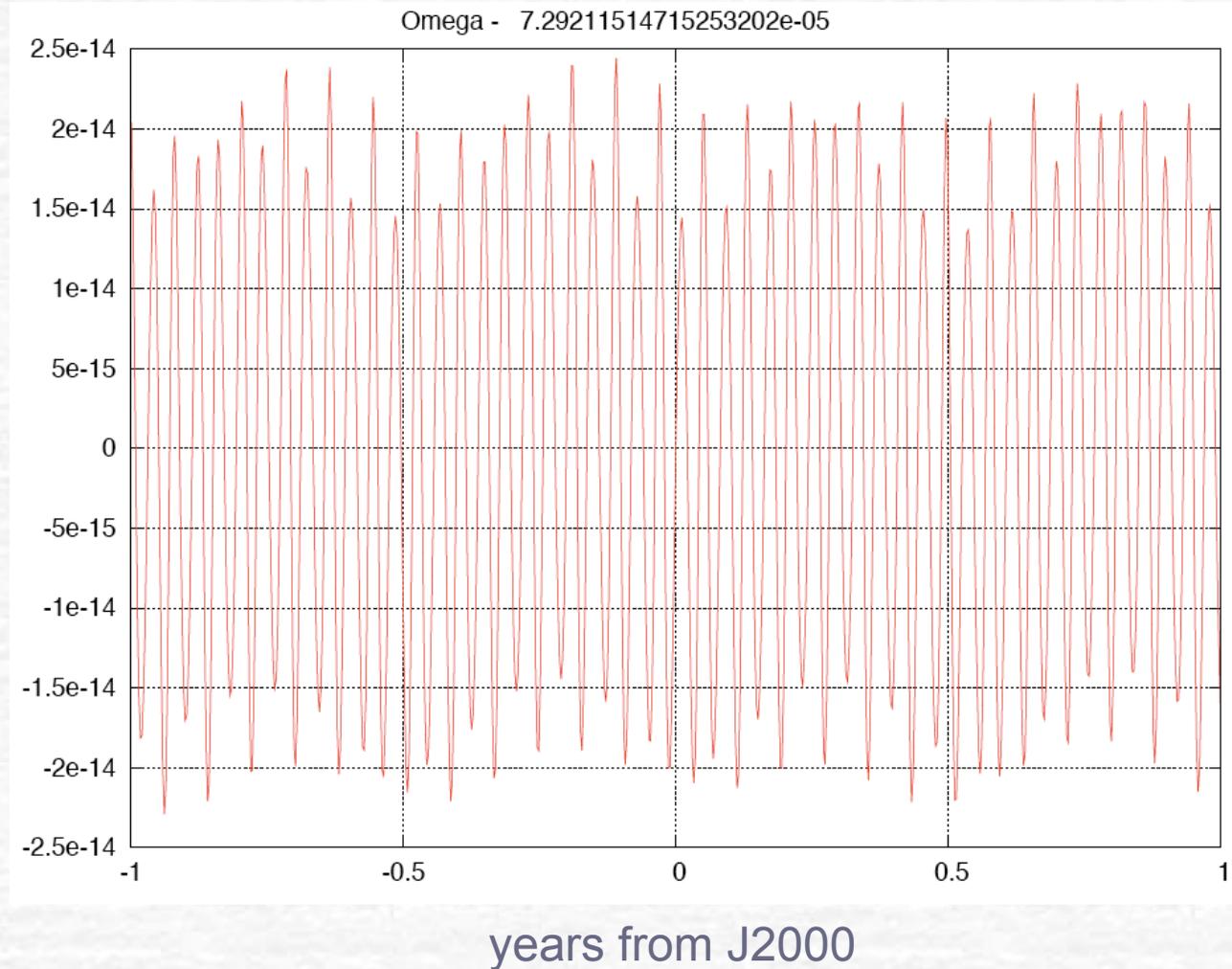


Effects on the LOD

6000 years with DE404:

$$\dot{\varphi} + \dot{\psi} \cos \omega = \Omega$$

rad/s



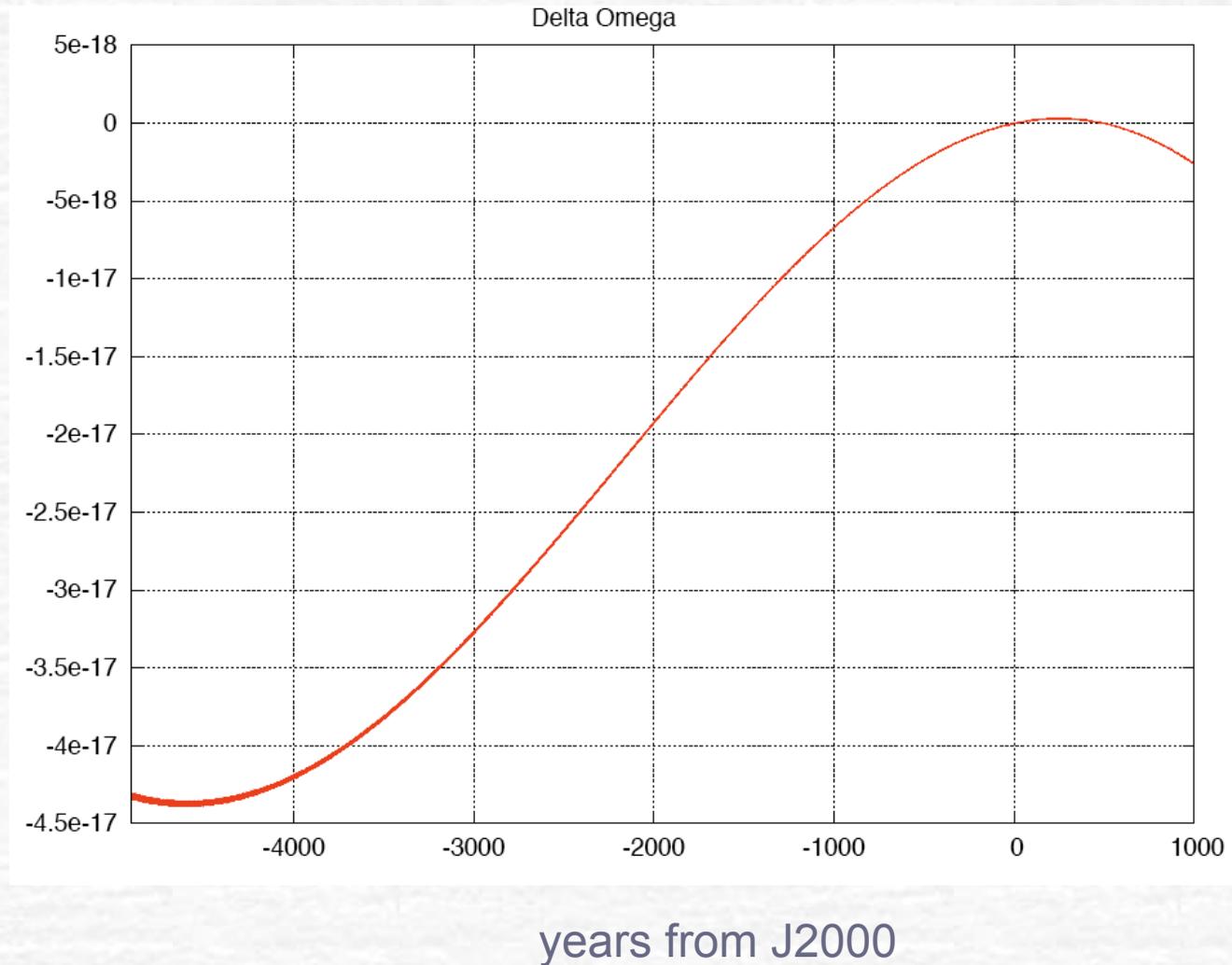
$$\left| \delta\Omega_N \right| < 450 \mu\text{as} / d \quad \Rightarrow \quad \Delta\text{LOD}_N < 30 \mu\text{s}$$

Effects on the LOD

6000 years with DE404:

$$\dot{\varphi} + \dot{\psi} \cos \omega = \Omega$$

rad/s



$$\left| \Omega_{pN} - \Omega_{Newt} \right| < 0.8 \mu as / d \quad \Rightarrow \quad \Delta LOD_{pN} < 0.06 \mu s$$