

POLARISATION I:
FOR THE LOVE OF STOKES

Jimi Green ATNF Radio Astronomy School 2009

Outline

- Why study Polarimetry and what is Polarisation?
- Poincare and his spheres
- Jones and his vectors (and matrices)
- Stokes and his parameters
- How do we measure Polarisation?
- Leakages and Mueller's Matrix
- Polarised beam effects
- Science with Polarisation:
 - ▣ Masers and Zeeman splitting



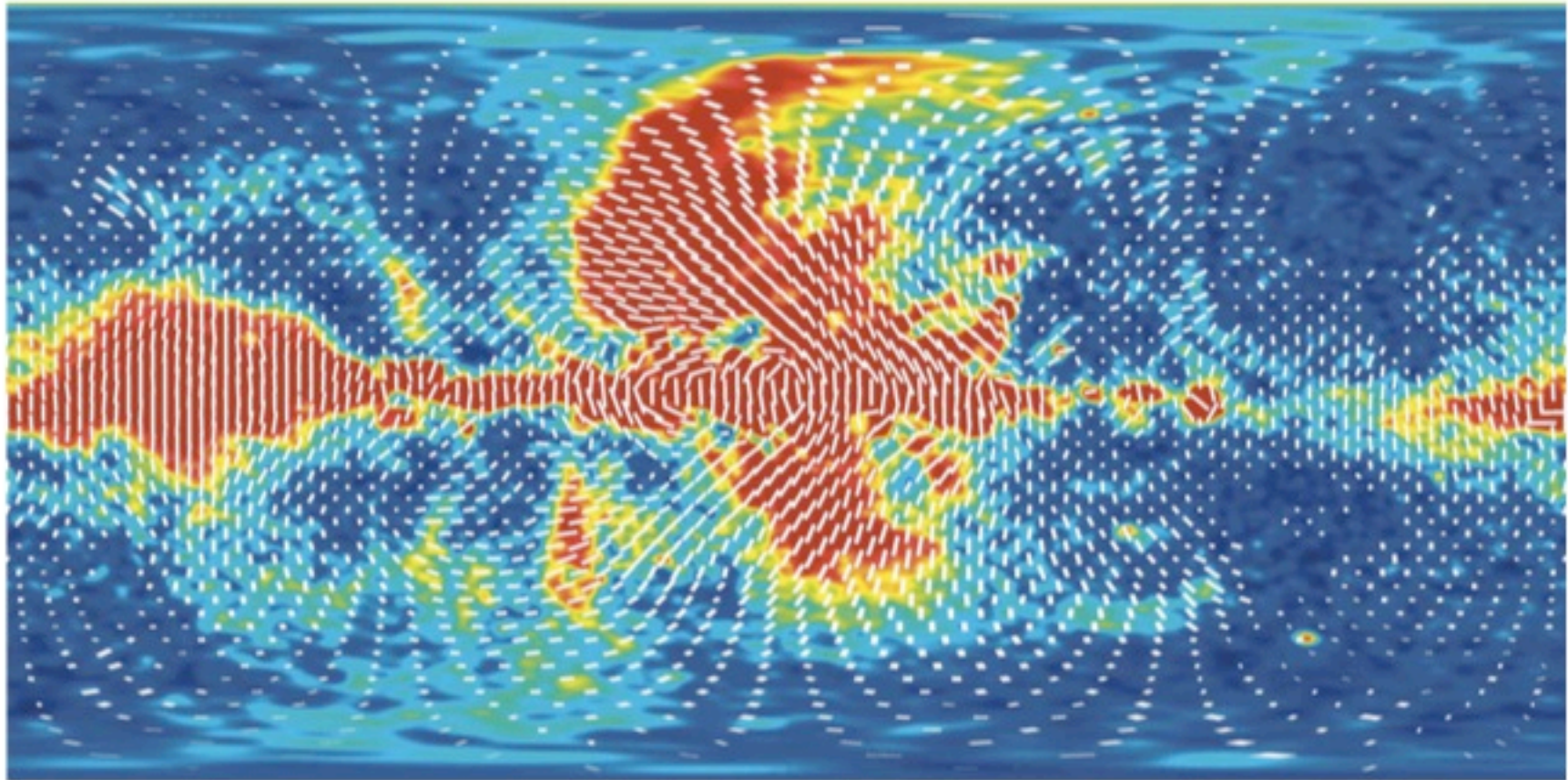
Why Polarimetry?

Why Polarimetry?

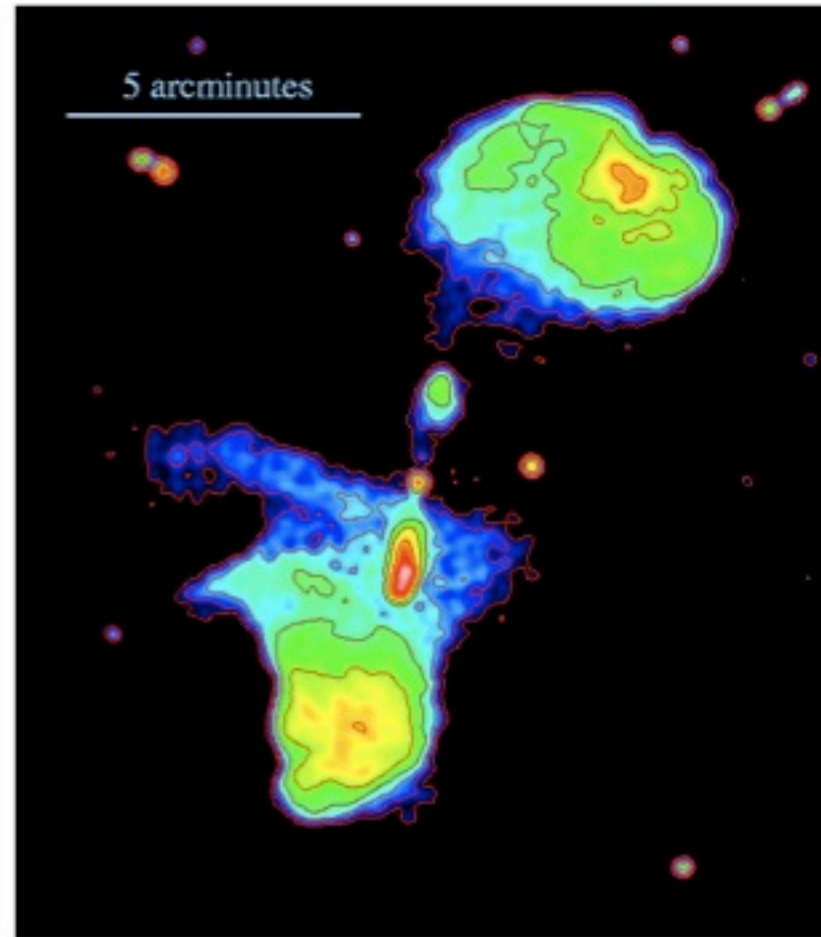


- Polarisation is fundamentally important to understanding the Universe
 - ▣ Provides insight into magnetic fields
- In optical astronomy, it's difficult to make polarimetric observations; in radio astronomy, they can be made easily, so why not use this to our advantage!
- (It's also very important to the Birds & the Bees, navigationally speaking, c.f. Rossel & Wehner, 1984)

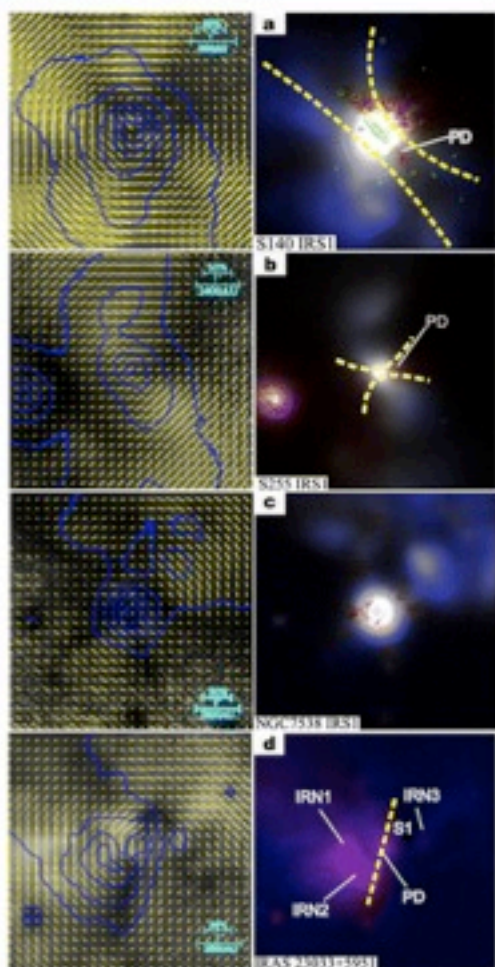
Why Polarimetry?



Why Polarimetry?



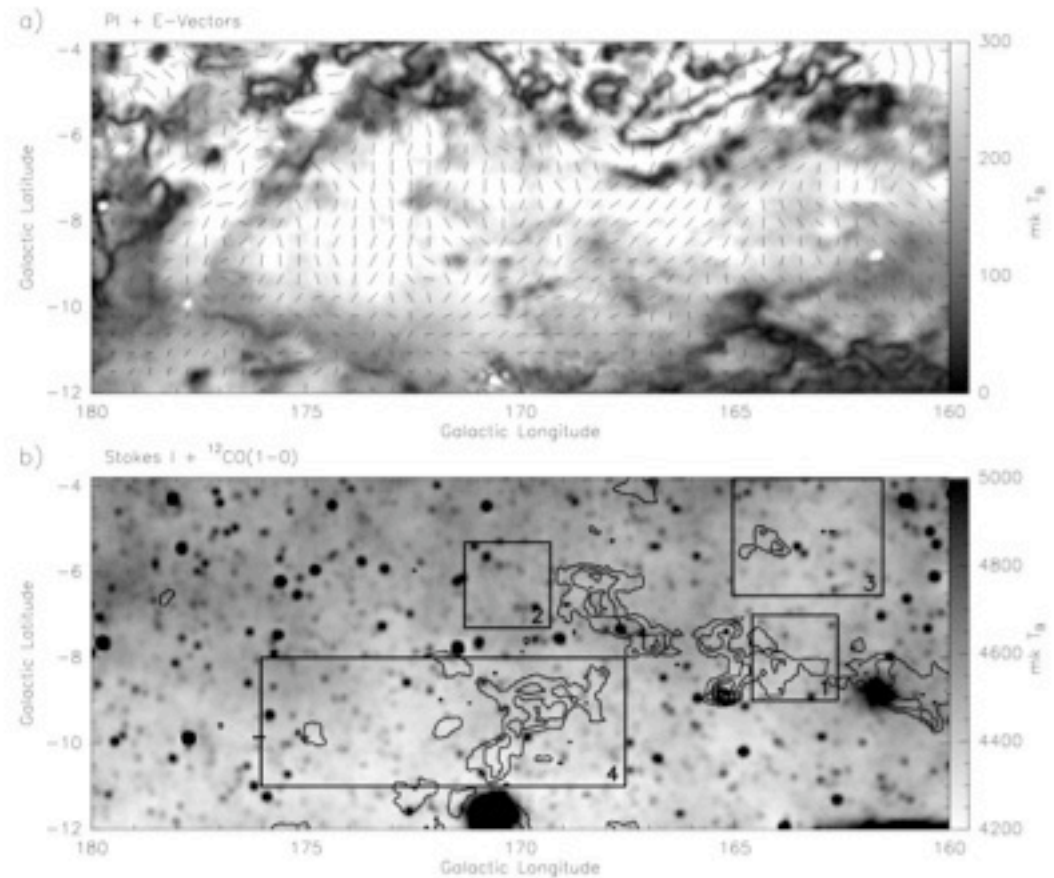
Why Polarimetry?



- *Left panels:* Polarization degree images overlaid by polarization vectors (yellow dashes) and total intensity contour (blue curves).
- *Right panels:* Pseudocolor images composed of pure brightness images (red) and polarized brightness images (blue).

Why Polarimetry

Faraday Screens have
only a polarised
structure, no standard
(Stokes I) emission!



Haverkorn et al.

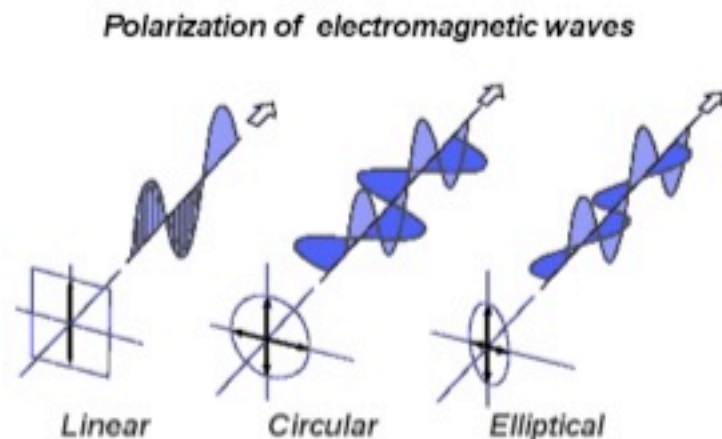
What is Polarisation?



What is Polarimetry & Polarisation?

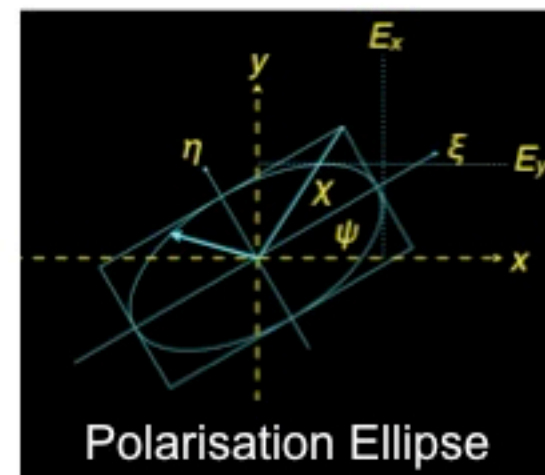
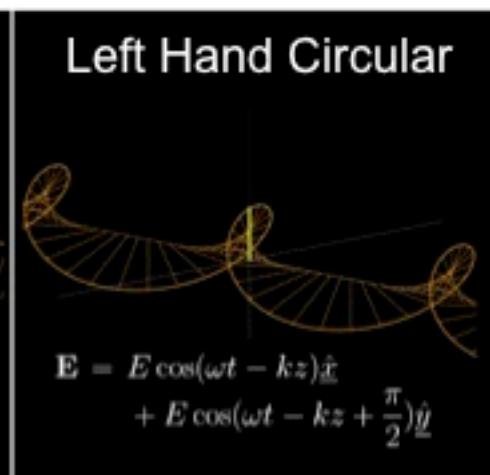
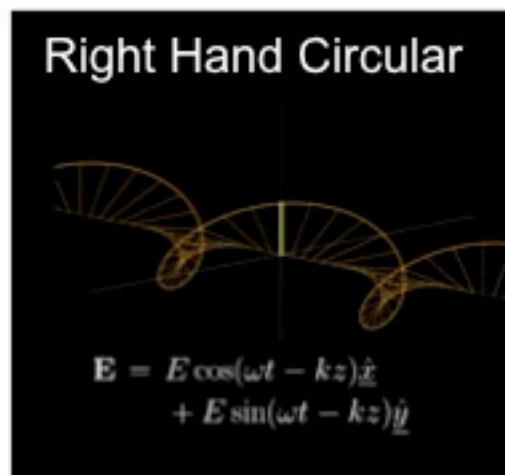
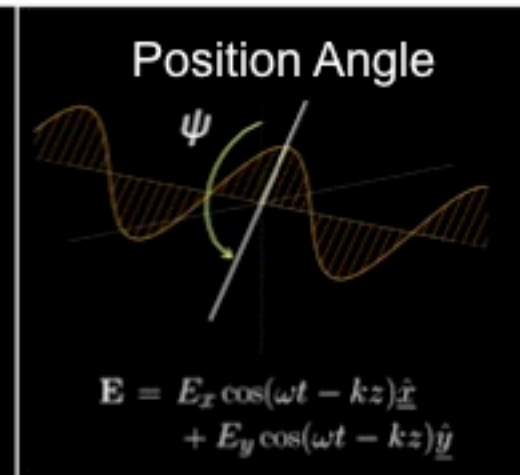
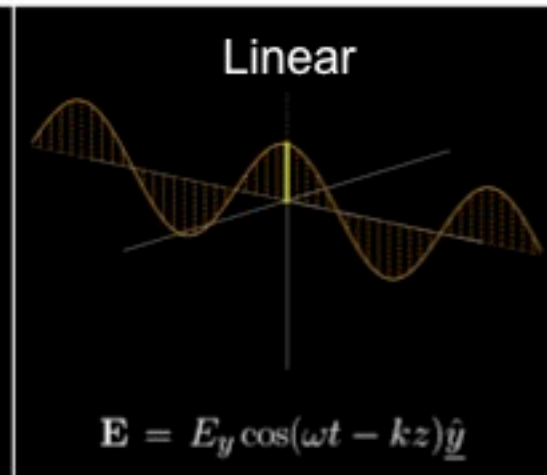
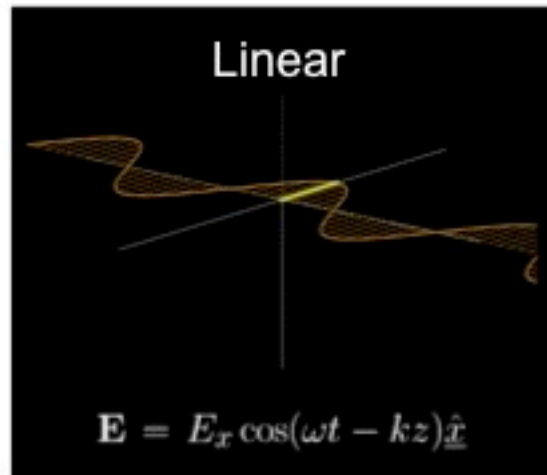
- Polarisation is the behaviour of the electric field with time.
- Natural radiation tends to be randomly polarised
 - ▣ The orientation of the electric field is completely random with respect to time.
- Astrophysical processes like synchrotron radiation can emit partially polarised emission, but *never* fully polarised.
- Interstellar matter can polarise random background emission or de-polarise polarised background emission.

What is Polarisation?



- Linear: orthogonal components in phase with constant ratio of strengths giving constant direction of electric vector.
- Circular: orthogonal components 90° out of phase with equal amplitudes – electric vector traces circle.

What is Polarisation?



What is Polarisation?

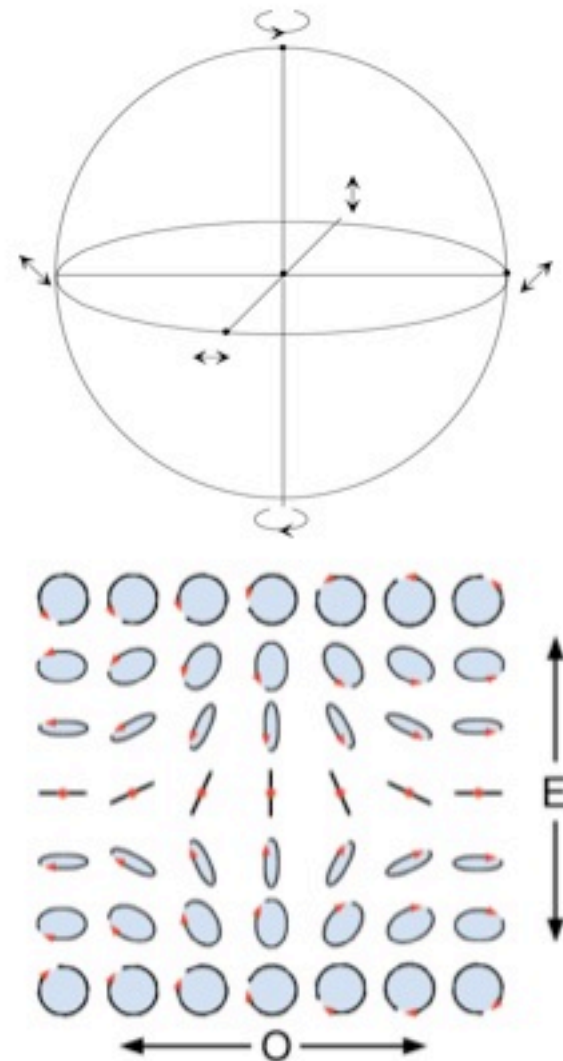
- Linearly polarised wave can be decomposed into two opposite handed circular waves
- Sum of two circular waves of unequal amplitude is elliptical.
- Sum of two orthogonal linears with a phase difference of between 0 and $\pi/2$ is also elliptical.



Jules Henri Poincaré

..and his sphere

- the spherical surface occupied by completely polarised states in the space of the vector
- Poles represent circular polarisations
 - ▣ Upper-hemisphere LHCP
 - ▣ Lower-hemisphere RHCP
- Equator represents linear polarisations with longitude representing tilt angle
- Latitude represents axial ratio



Now for the maths





Robert Clarke Jones

..and his vectors

- Jones calculus is a matrix-based means of relating observed to incident fields.
- Vectors describe incident radiation and matrices the response of the instrument.
- The Jones Vector: $\begin{pmatrix} E_x(t) \\ E_y(t) \end{pmatrix}$
- Examples:
 - ▣ Linearly (x-direction) polarised wave: $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$
 - ▣ Left-Hand Circularly polarised wave: $\frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}$

Robert Clarke Jones



..and his matrices

- Effect of instrument described by 2x2 matrix:

$$\begin{pmatrix} E_x \\ E_y \end{pmatrix}_o = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} E_x \\ E_y \end{pmatrix}_i$$

- Simple Examples:

- Linear polariser: $\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$

- Left-Hand Circular polariser: $\frac{1}{2} \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix}$

- In practice matrix elements complex.
- Important: Only applicable to **completely** polarised waves.



Sir George Gabriel Stokes

..and his parameters

- Defined by George in 1852
- Adopted for astronomy by Chandrasehkar in 1947.
- Can be used for partially polarised radiation.
- Not a vector quantity! Deals with power instead of electric field amplitudes.
- The correlator can produce ALL Stokes parameters simultaneously (not so easy in optical astronomy!)

Stokes Parameters

- I – total intensity and sum of any two orthogonal polarisations
- Q and U – completely specify linear polarisation
- V – completely specifies circular polarisation

$$I = E_{0x}^2 + E_{0y}^2$$

$$Q = E_{0x}^2 - E_{0y}^2$$

$$U = 2E_{0x}E_{0y} \cos \delta$$

$$V = 2E_{0x}E_{0y} \sin \delta$$

$$I = \langle E_x E_x^* \rangle + \langle E_y E_y^* \rangle$$

$$Q = \langle E_x E_x^* \rangle - \langle E_y E_y^* \rangle$$

$$U = \langle E_x E_y^* \rangle + \langle E_x^* E_y \rangle$$

$$V = i \left(\langle E_x E_y^* \rangle - \langle E_x^* E_y \rangle \right)$$

(For linear feeds)

Fractional Polarisations

- The total linearly polarised intensity is defined as:

$$P = \sqrt{U^2 + Q^2} \quad \text{[for native linear feed]}$$

- A linearly polarised source will have an intrinsic position angle on the sky that is given by:

$$\Theta = \frac{1}{2} \tan^{-1} \left(\frac{U}{Q} \right) \quad \text{[for native linear feed]}$$

- The circular polarisation will be just Stokes V.
- Stokes parameters often presented as percentages of the total intensity.
- Since radio sources are *never* fully polarised, then the fractional linear and circular polarisation will always be < 1

How do we measure it?



How do we measure it?

- Stokes parameters are the auto-correlation & cross-correlation products returned from the correlator, but input to the correlator can come from different feed types.
- Feeds normally designed to approximate pure linear or circular (known as 'native linear' or 'native circular')
 - ▣ Linear Feeds – intrinsically accurate & provide true linear response.
 - ▣ Circular Feeds – less accurate & frequency dependent response.

How do we measure it?

- Output of native linear feed is E_x and E_y field voltages, so:
 - I from $XX+YY$
 - Q from $XX-YY$
- Native circular adds 90° phase to X, so:
 - I from $XX+YY$
 - V from $XX-YY$

Stokes Parameters

- For circular feeds Q and V swap round..

Linear

$$XX = I + Q$$

$$YY = I - Q$$

$$XY = U + iV$$

$$YX = U - iV$$

Circular

$$RR = I + V$$

$$LL = I - V$$

$$RL = Q + iU$$

$$LR = Q - iU$$

But is it really that simple?

- Do we just plug in our computer and get $\{I, Q, U, V\}$ out of the correlator?
- No, there are leakages!
 - ▣ The total intensity can leak into the polarised components (I into $\{Q, U, V\}$).
 - ▣ The linear polarisation can leak into the circular ($\{Q, U\}$ into V).
 - ▣ ... and all combinations and permutations are allowed!
- Without correcting for leakage, you're not going to get proper Stokes parameters!



Hans Mueller

..and his matrix

- The leakage of each polarisation into the other can be measured and quantified in a 4x4 matrix first proposed by Mueller in 1943.

$$M = \begin{bmatrix} m_{II} & m_{IQ} & m_{IU} & m_{IV} \\ m_{QI} & m_{QQ} & m_{QU} & m_{QV} \\ m_{UI} & m_{UQ} & m_{UU} & m_{iI} \\ m_{VI} & m_{VQ} & m_{VU} & m_{VV} \end{bmatrix}$$

The Mueller Matrix

Correlator Output

Incoming Radiation

$$\begin{bmatrix} XX + YY \\ XX - YY \\ XY \\ YX \end{bmatrix} = M \cdot \begin{bmatrix} I \\ Q \\ U \\ V \end{bmatrix}$$

Example (simple) Mueller Matrices

- If feeds were perfect:

- Dual linear feed: M is unitary

- Dual linear feed rotated 45° : Q and U interchange and sign change for rotation:

$$M = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

- A dual linear feed rotated 90° : signs of Q and U reversed:

$$M = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

- As Alt-Az telescope tracks source, feed rotates on sky by the parallactic angle (PA):

$$M_{sky} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos 2PA & \sin 2PA & 0 \\ 0 & -\sin 2PA & \cos 2PA & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

The more general Mueller Matrix

- For a (realistic) dual linear feed:

$$M = \begin{bmatrix} 1 & \left(-2\varepsilon \sin\phi \sin 2\alpha + \frac{\Delta G}{2} \cos 2\alpha\right) & 2\varepsilon \cos\phi & \left(2\varepsilon \sin\phi \cos 2\alpha + \frac{\Delta G}{2} \sin 2\alpha\right) \\ \frac{\Delta G}{2} & \cos 2\alpha & 0 & \sin 2\alpha \\ 2\varepsilon \cos(\phi + \varphi) & \sin 2\alpha \sin\varphi & \cos\varphi & -\cos 2\alpha \sin\varphi \\ 2\varepsilon \sin(\phi + \varphi) & -\sin 2\alpha \cos\varphi & \sin\varphi & \cos 2\alpha \cos\varphi \end{bmatrix}$$

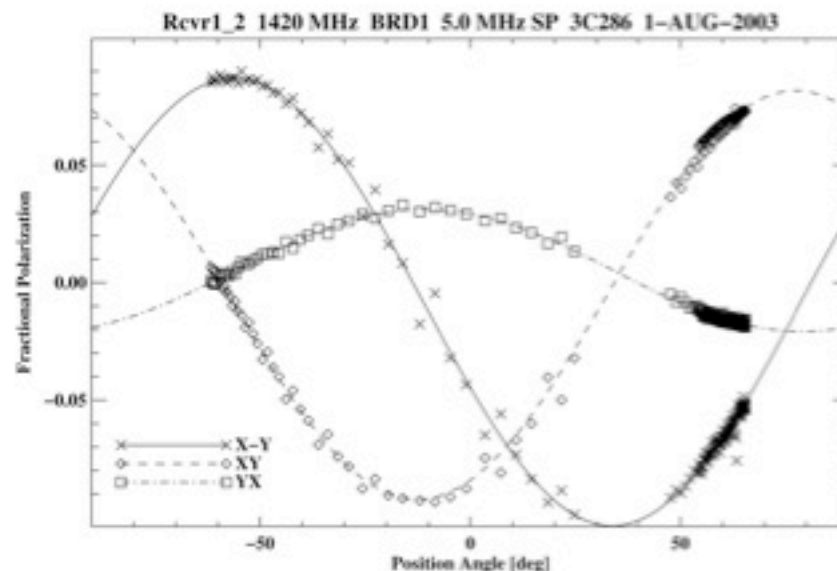
- The Mueller matrix has 16 elements, but **ONLY 7 INDEPENDENT PARAMETERS**. The matrix elements are not all independent.

Calculating the Mueller Matrix

- For a perfect system, as we track a polarised source across the sky the parallactic angle changes and this should produce:
 - For XX-YY: $\cos 2(\text{PA}_{\text{az}} + \text{PA}_{\text{src}})$, centred at zero.
 - For XY: $\sin 2(\text{PA}_{\text{az}} + \text{PA}_{\text{src}})$, centred at zero
 - For YX: zero (most sources have zero circular polarisation)

Calculating the Mueller Matrix

- But, what we find is:



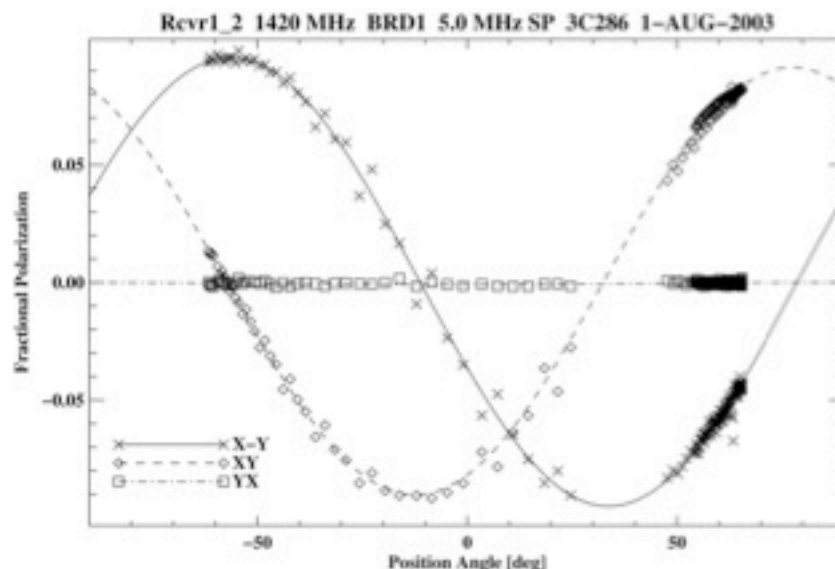
DELTA Γ = -0.017 ± 0.006
PSI = -16.7 ± 1.8
ALPHA = $+0.4 \pm 0.9$
EPSILON = $+0.004 \pm 0.001$
PHI = $+152.3 \pm 23.0$
Q_{SAC} = -0.037 ± 0.002
U_{SAC} = -0.085 ± 0.002
POI_{SAC} = $+0.093 \pm 0.002$
PA_{SAC} (**UNCORRECTED FOR M_{EXTRO}**) = -56.81 ± 0.64
NR GOOD POINTS: X-Y = 149 XY = 158 YX = 160 / 160

Mueller Matrix:

1.0000	-0.0083	-0.0065	0.0033
-0.0083	0.9999	0.0001	0.0145
-0.0053	-0.0042	0.9581	0.2865
0.0052	-0.0139	-0.2865	0.9580

Calculating the Mueller Matrix

- Which enables the matrix to be calculated and the observations corrected to give what we expect:



DELTA θ = 0.000 ± 0.006
PSI = $+0.3 \pm 1.8$
ALPHA = -0.0 ± 0.9
EPSILON = $+0.000 \pm 0.001$
PHI = -39.8 ± 235.4
Q_{SAC} = -0.037 ± 0.002
U_{SAC} = -0.085 ± 0.002
POL_{SAC} = $+0.093 \pm 0.002$
PA_{SAC} (**UNCORRECTED FOR M_{SYSTEM}**) = -56.82 ± 0.64
NR GOOD POINTS: X-Y = 149 XY = 158 YX = 158 / 160

Mueller Matrix:

1.0000	0.0002	0.0006	-0.0005
0.0002	1.0000	0.0000	-0.0012
0.0006	-0.0000	1.0000	-0.0044
-0.0005	0.0012	0.0044	1.0000

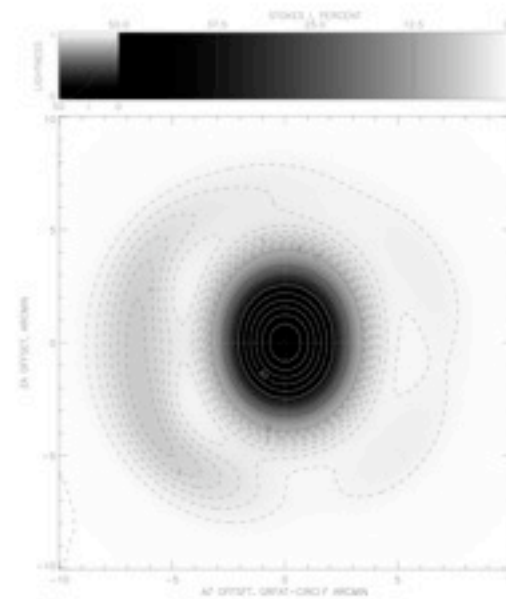
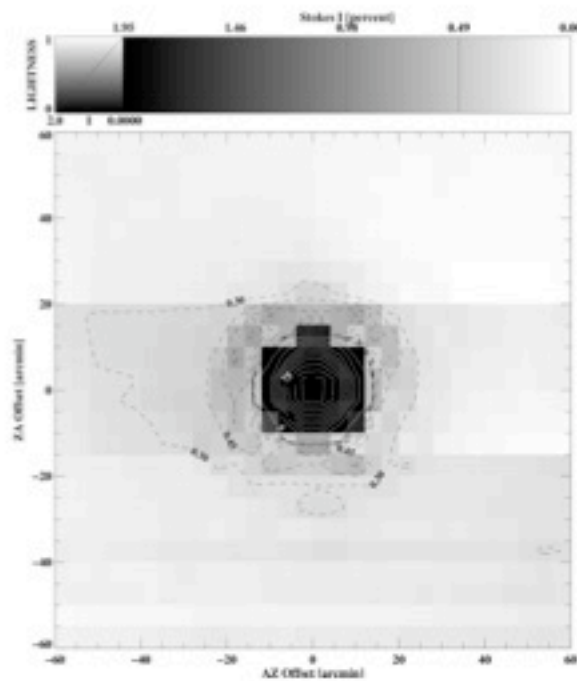


Beam Effects

Beam Effects

- For point sources, all of the previous is fine.
- What if the source you're looking at is extended compared to the telescope beam?
- There are instrumental beam effects that can confuse the measurement of extended polarised signals. They are...
 - ▣ Squint
 - ▣ Squash

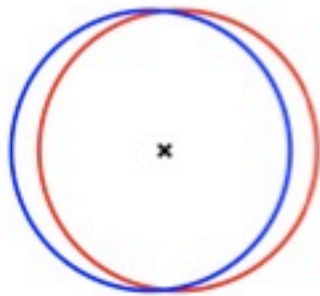
Stokes I response



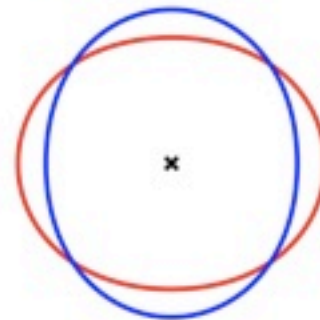
Heiles et al. 2001

Beam Squint & Squash

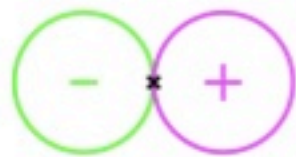
BEAM SQUINT
RHCP
LHCP



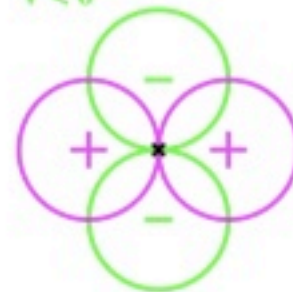
BEAM SQUASH
RHCP
LHCP



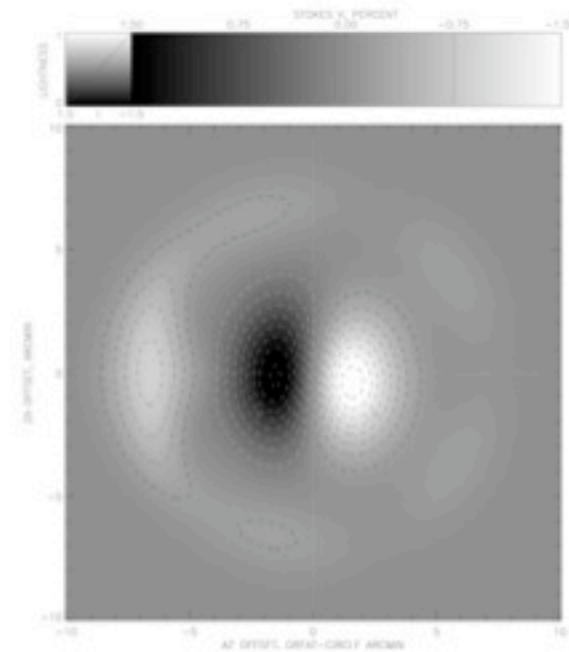
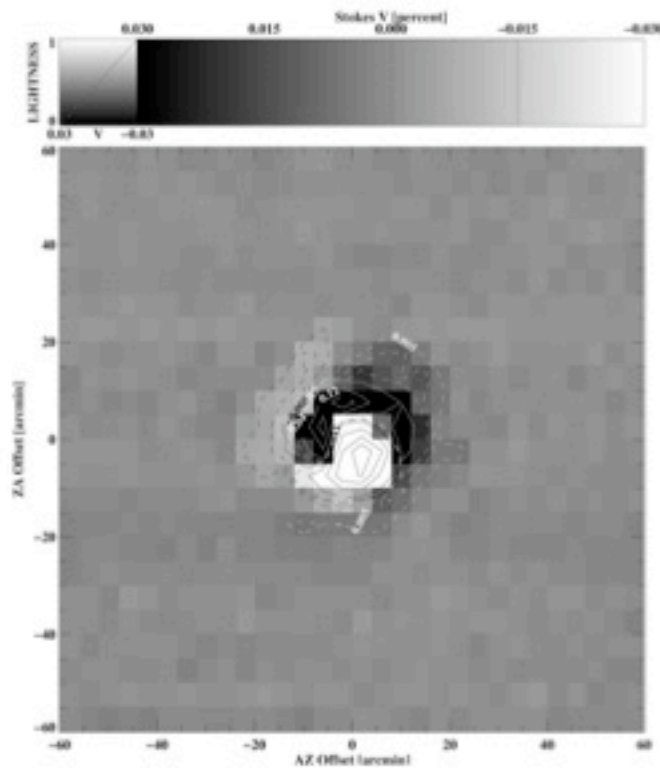
$V = \text{RHCP} - \text{LHCP}$
 $V > 0$
 $V < 0$



$V = \text{RHCP} - \text{LHCP}$
 $V > 0$
 $V < 0$

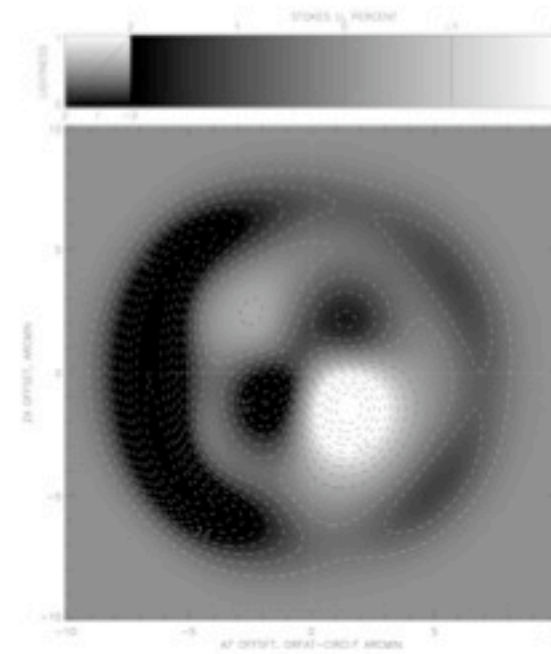
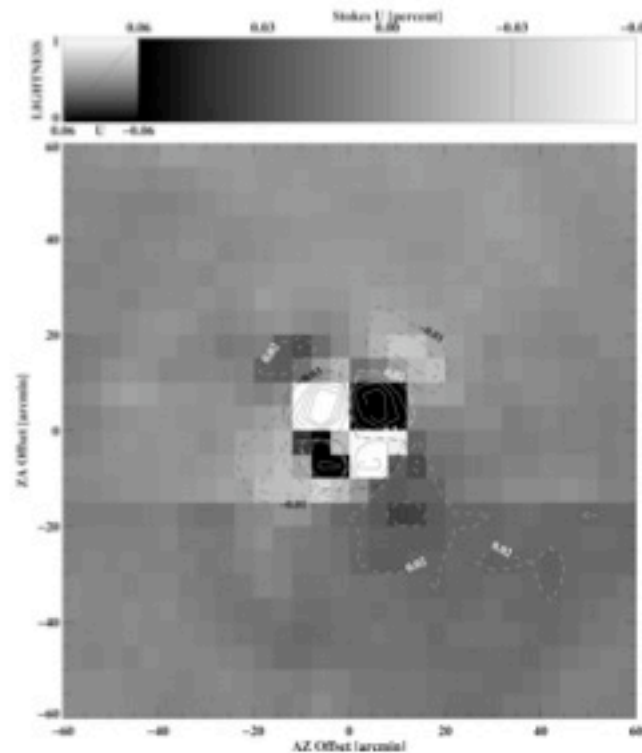


Squint in action



Heiles et al. 2001

Squash in action

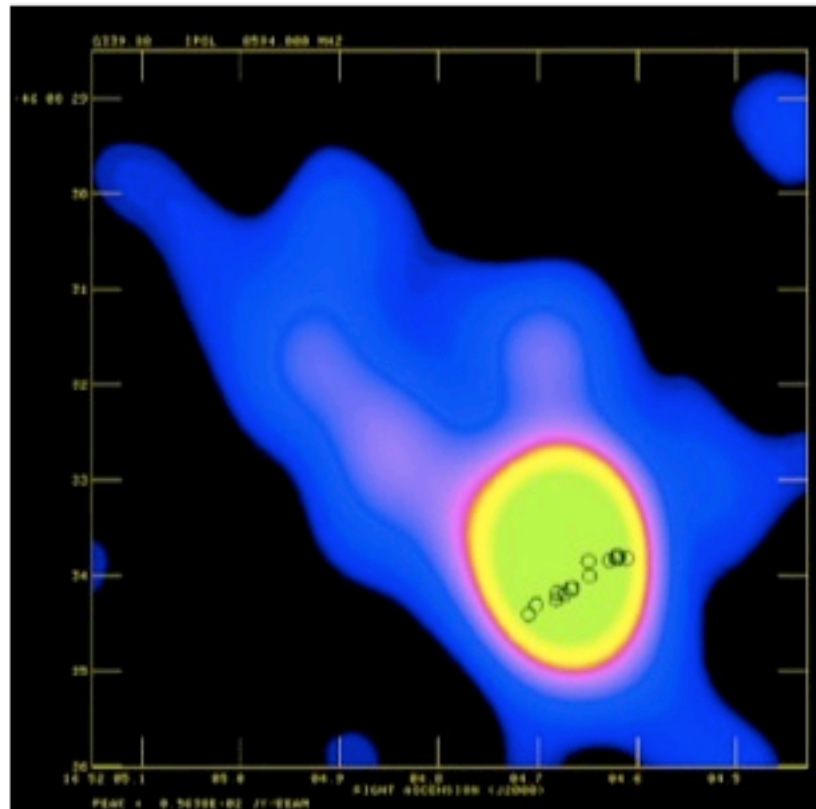


Heiles et al. 2001

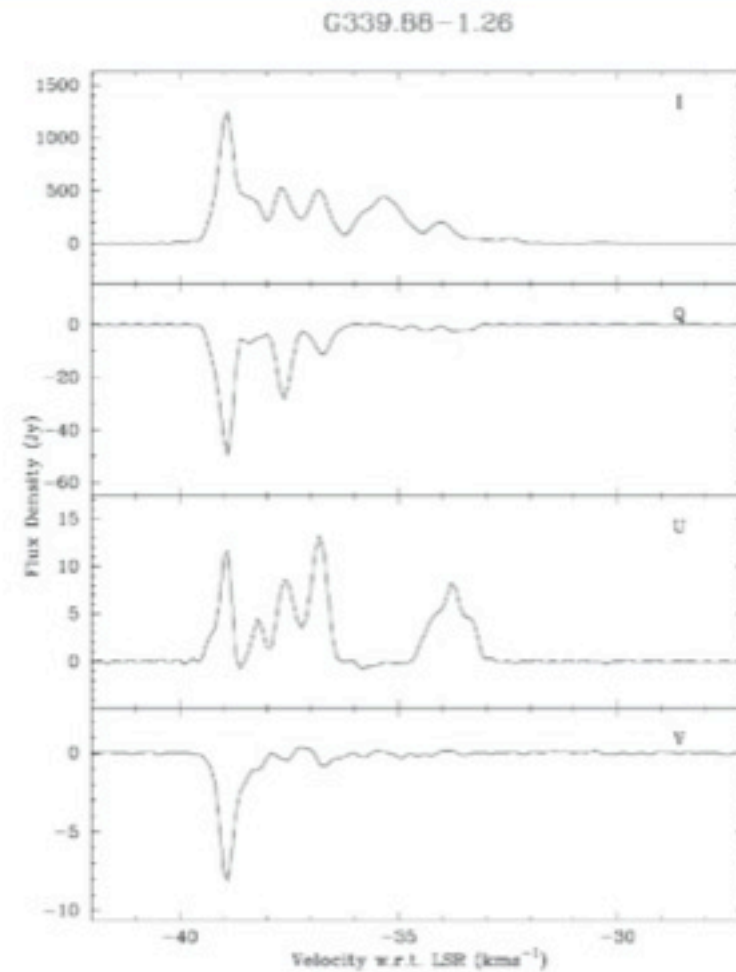
Here comes the science



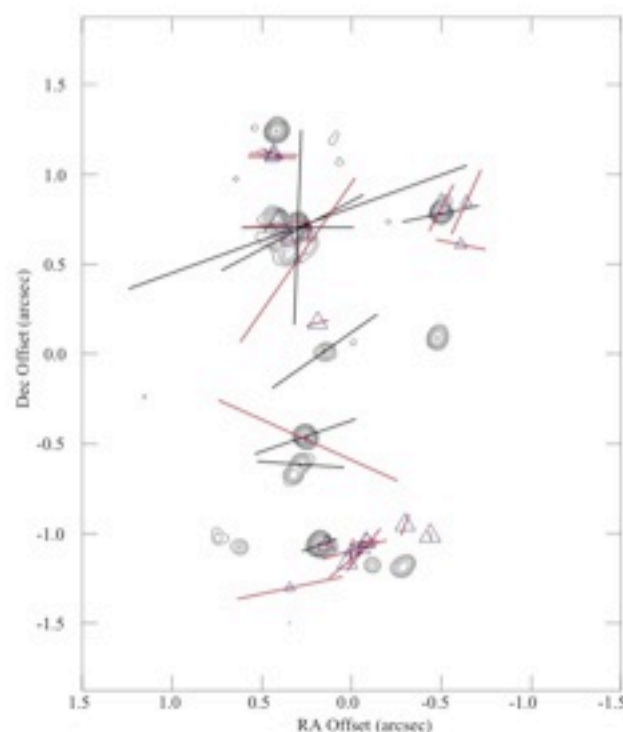
Polarisation of Masers



S. Ellingsen



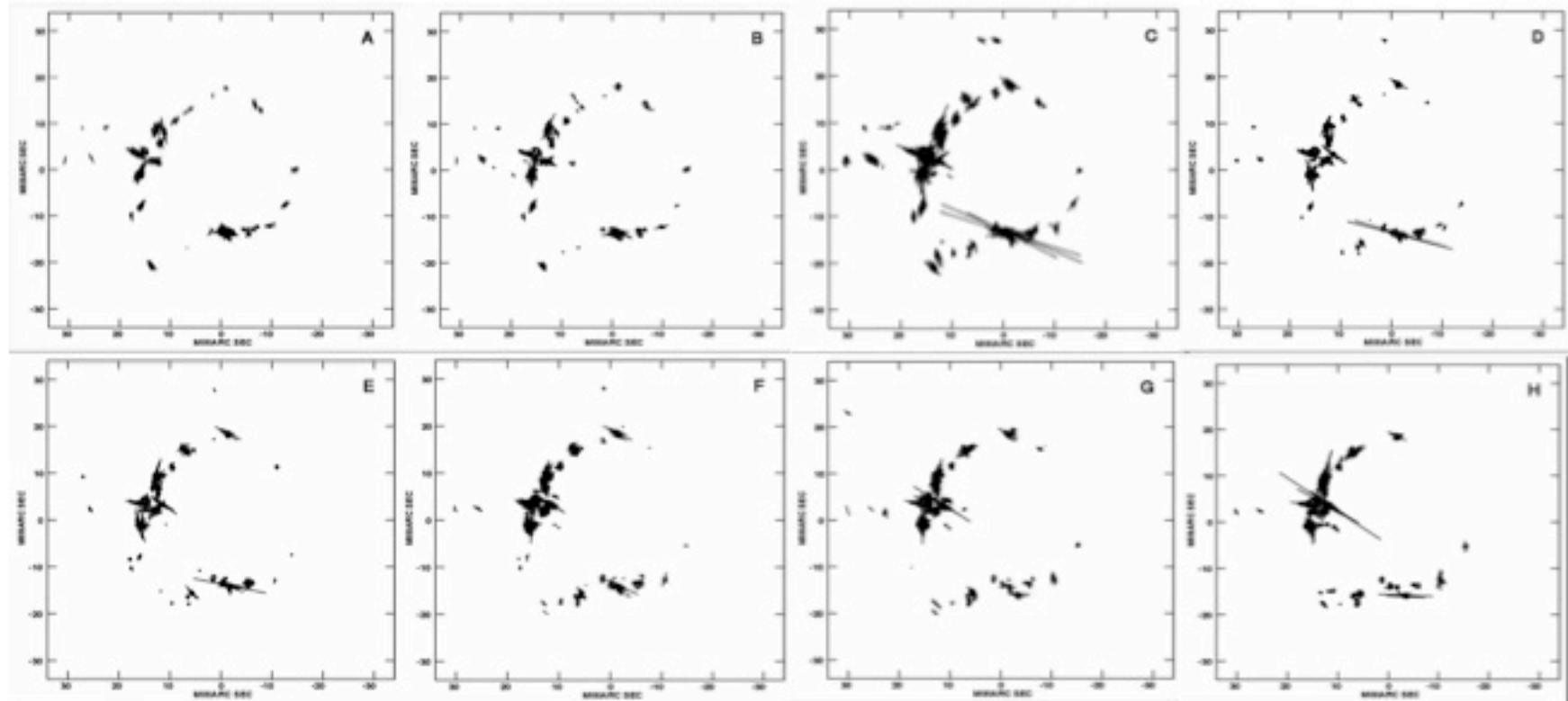
Polarisation of Masers



W3(OH) Harvey-Smith & Cohen 2006, Vlemmings et al 2006.

- Filamentary maser structure.
- Linear polarisation up to 8%.
- Polarisation angles indicate north-south structure, and are consistent with OH.

Linear Polarisation of Masers



Kemball et al. 2009

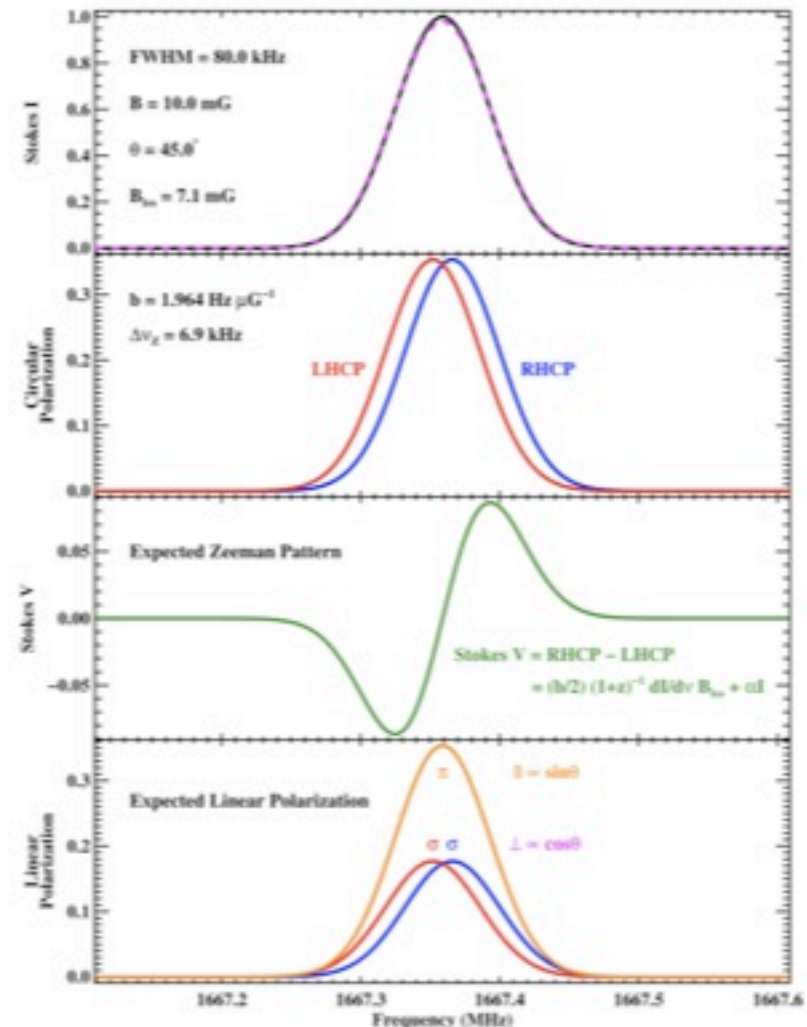
Make like a banana:

Zeeman Splitting

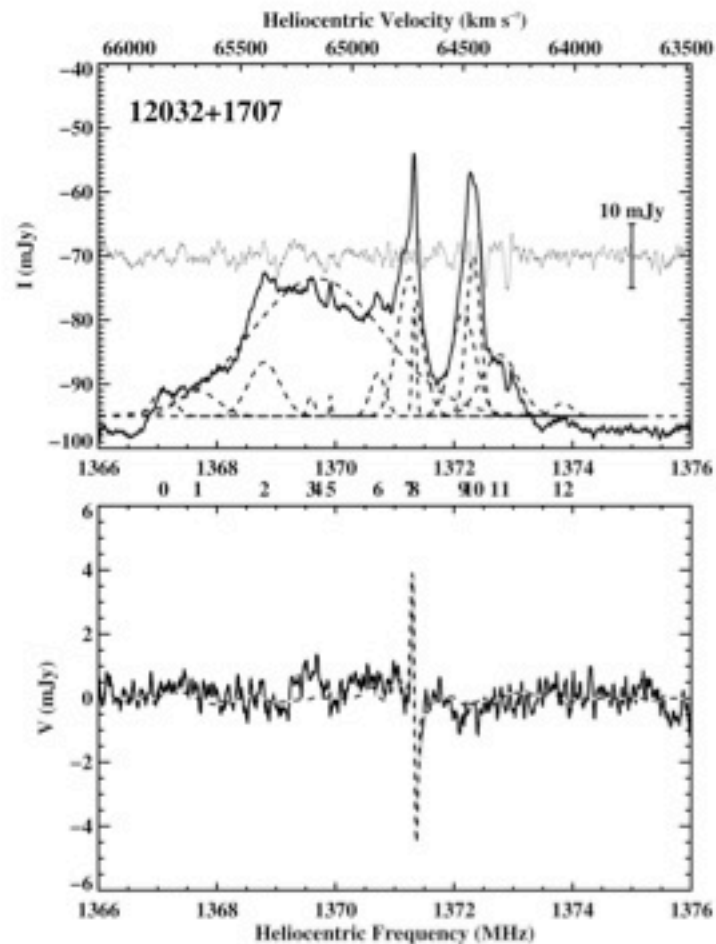


- Atoms & molecules in net magnetic moment will have their energy levels split in the presence of a magnetic field.
- Detected through frequency shift between right and left circularly polarised emission

$$V = RHCP - LHCP \propto B_{los}$$



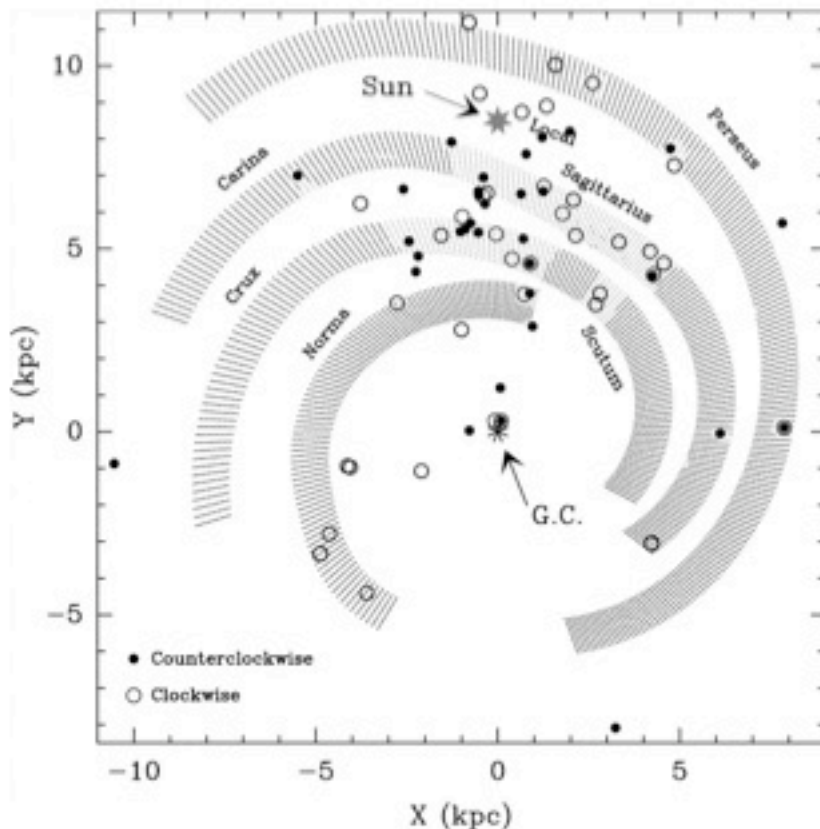
Zeeman Splitting of Masers



- OH Gigamaser
(luminosity greater than 10^4 solar luminosities)



Zeeman Splitting of Masers



Fish et al 2003

- Line-of-sight magnetic field directions deduced from OH maser Zeeman splitting.
- 74 star-forming regions:
 - ▣ 41 with an overall magnetic field oriented in a clockwise sense.
 - ▣ 33 with field oriented counterclockwise as viewed from above the Galactic center.
- Field consistency within 2-kpc of Sun.



In summary..

Summary

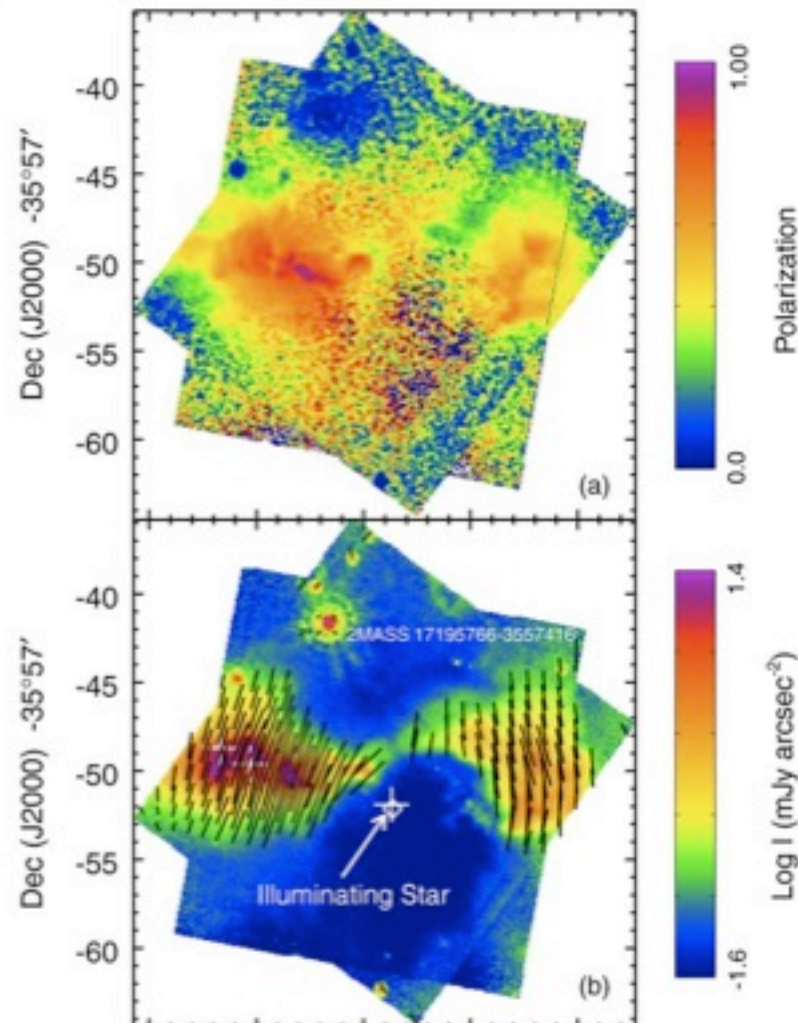
- Polarisation in radio astronomy very important to improving our knowledge & understanding.
- Can describe polarisation with the Polarisation Ellipse and the Poincare Sphere.
- Dr. Jones offers a vector representation for ideal cases of completely polarised emission.
- Mueller and his matrices are the best option for real situations.
- There are Linear and Circular feed types, must account for which you are using.
- **Understanding the polarisation properties of your dish is fundamental to successful observations!**
- (Masers offer exciting science opportunities!)

Useful References

- Heiles, C. 'A Heuristic Introduction to Radioastronomical Polarisation' (2002) ASP 278
- Tinbergen, J. 'Astronomical Polarimetry' (1996), Cambridge University Press (Cambridge, UK)
- Stutzman, W. 'Polarisation in Electromagnetic Systems' (1993), Artech House (Norwood, MA, USA)
- Radhakrishnan. Polarisation. URSI proceedings (1990) pp. 34
- Hamaker et al. Understanding radio polarimetry. I. Mathematical foundations. Astronomy and Astrophysics Supplement (1996) vol. 117 pp. 137
- Born and Wolf: 'Principle of Optics', Chapters 1 and 10



Why Polarimetry?



- *HST* NICMOS image of massive young stellar object NGC 6334 V.
- (a) Fractional polarization.
- (b) Log intensity with polarization vectors.



Michael Faraday

..and his rotations

- Magnetised plasma is birefringent – the refractive indices for the two circular modes are different due to the parallel component of the magnetic field (and dependent on electron density and frequency of radiation).
- Hence phases of two modes changes along the propagation path and so does the position angle of the resultant linearly polarised radiation.